

## Deformation and orientation effects in the binary symmetric decay of $^{20,21,22}\text{Ne}^*$

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### Introduction

We have extended the study of binary symmetric decay (BSD) of extremely light mass compound systems  $^{20,21,22}\text{Ne}^*$  formed in  $^{10,11}\text{B}+^{10,11}\text{B}$  reactions at  $E_{\text{lab}} = 48$  MeV, to explore the role of deformations and orientations, using the Dynamical Cluster decay Model (DCM) of Gupta and collaborators [1–3]. The results are compared with purely spherical consideration of nuclei to bring out the effects of oriented nuclei as well as with that of our earlier study [2] having spherical consideration within the description of Süsmann central radii  $C_t = C_1 + C_2$ , to assimilate the deformation effects, where  $C_i = R_i - b^2/R_i$  (in fm) with  $R_i = [1.28A_i^{1/3} - 0.76 + 0.8A_i^{1/3}][1 + 0.0007T^2]$  fm and surface thickness parameter  $b=0.99$  fm.

Our earlier studies [2] for the BSD of  $^{20,21,22}\text{Ne}^*$  systems reveals fusion-fission  $\sigma_{ff}$  and deep inelastic orbiting  $\sigma_{DIO}$  contributions from compound nucleus CN and non-compound nucleus nCN processes, respectively, in the total DCM calculated BSD cross sections  $\sigma_{BSD}^{DCM}$ . On the basis of preformation probabilities (Fig. 1(a-c)), we noticed that the BSD of  $^{20}\text{Ne}^*$  is highly favored whereas for  $^{21}\text{Ne}^*$  there is strong competition from neighbouring fragments, whereas for  $^{22}\text{Ne}^*$  BSD is least favored. It pointed that nuclear structure is playing significant role in BSD of these systems and further emphasized that the process of collective clusterization is playing very strong role within the DCM. In the present work [3], we find that with inclusion of quadruple deformations and “hot compact” orienta-

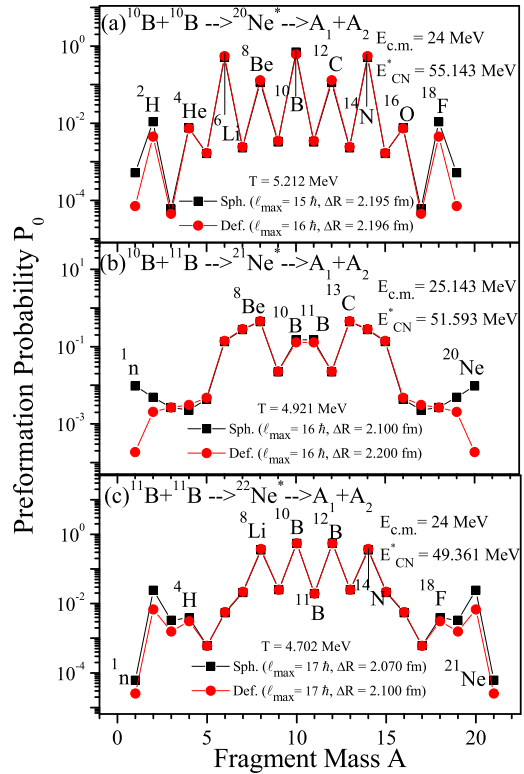


FIG. 1: Preformation Probability  $P_0$  as a function of fragment mass  $A$  for the decay of a)  $^{20}\text{Ne}^*$ , b)  $^{21}\text{Ne}^*$ , and c)  $^{22}\text{Ne}^*$ , at respective  $\ell_c$  for spherical and oriented nuclei considerations.

tions of nuclei  $\sigma_{ff}$  increases in comparison to the case of spherical considerations of nuclei.

### Methodology

The decay of hot and rotating compound nucleus is studied within framework of dynamical cluster decay model (DCM), which is worked out in terms of collective coordinates of mass asymmetry  $\eta = (A_T - A_P)/(A_T + A_P)$

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TABLE I: The BSD cross-sections calculated with  $\ell$  summed upto  $\ell_c$  for the decay of compound systems  $^{20}\text{Ne}^*$ ,  $^{21}\text{Ne}^*$ , and  $^{22}\text{Ne}^*$ , using the DCM, alongwith their comparison with the experimental data [4].

System	Parameter	$\ell_c$ ( $\hbar$ )	$\Delta R_{ff}$ (fm)	$\Delta R_{DIO}$ (fm)	$\sigma_{BSD}$ (mb)			
					DCM			Expt.
					<i>ff</i>	<i>DIO</i>	Total	
$^{20}\text{Ne}^*$		15	2.195	1.683	236.46	33.30	269.76	$\sim 270$
$^{21}\text{Ne}^*$	$R_t + \Delta R$ ( <i>Sph.</i> )	16	2.100	1.574	102.60	27.37	129.97	$< 130$
$^{22}\text{Ne}^*$		17	2.070	1.605	8.52	60.48	69.0	$< 70$
$^{20}\text{Ne}^*$	$C_t + \Delta R$	12	1.490	0.880	195.27	105.21	300.48	$\sim 270$
$^{21}\text{Ne}^*$		13	1.390	0.725	65.723	93.02	158.74	$< 130$
$^{22}\text{Ne}^*$		14	1.490	0.690	8.677	88.62	97.297	$< 70$
$^{20}\text{Ne}^*$	$R_t + \Delta R$ ( <i>Def.</i> )	16	2.196	1.333	267.78	2.11	269.89	$\sim 270$
$^{21}\text{Ne}^*$		16	2.200	1.632	103.62	26.26	129.88	$< 130$
$^{22}\text{Ne}^*$		17	2.100	1.643	8.68	59.34	68.02	$< 70$

and relative separation (R) with effects of temperature, deformation and orientation duly incorporated in it. In terms of these collective coordinates, using the  $\ell$ - partial waves, the decay cross-section is defined as

$$\sigma = \frac{\pi}{k^2} \sum_{\ell=0}^{\ell_c} (2\ell + 1) P_0 P; \quad k = \sqrt{\frac{2\mu E_{c.m.}}{\hbar^2}} \quad (1)$$

where preformation probability  $P_0$  refers to  $\eta$  motion and is given by sol. of stationary Schrodinger eq. in  $\eta$ . Penetrability P refers to R motion and is calculated using WKB approximation,  $\mu$  is the reduced mass and  $\ell_c$ , the critical angular momentum.

### Calculations and discussions

Fig. 1 (a-c) gives preformation probability  $P_0$  as a function of fragment mass A for  $^{20}\text{Ne}^*$ ,  $^{21}\text{Ne}^*$  and  $^{22}\text{Ne}^*$ , respectively. The comparison between considerations of purely spherical and oriented nuclei gives similar results as it is evident from Fig. 1 as well as Table. 1. As we also found in our earlier study [2], the BSD for  $^{20}\text{Ne}^*$  is highly favored followed by  $^{21}\text{Ne}^*$  and least favored in case of  $^{22}\text{Ne}^*$ . Hence it is very much evident that the nuclear structure effects are playing an important role, within DCM, via preformation probability of different fragments prior to their decay. Table. 1 presents the comparative analysis of total  $\sigma_{BSD}^{DCM}$  with the  $\sigma_{BSD}^{Expt.}$  data for pure

spherical, spherical ( $C_t + \Delta R$ ) and oriented nuclei considerations. The contribution of  $\sigma_{DIO}$  is calculated by considering the  $P_0 = 1$  for incoming channel i.e. the incoming nuclei are considered not to lose their identity. We note that for oriented nuclei consideration, the  $\sigma_{ff}$  contribution is more in comparison to the spherical consideration of fragments particularly for  $^{20}\text{Ne}^*$  system. The results with the considerations of purely spherical and oriented nuclei show good agreement with the experimental data [4].

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