

Study of s-state Ξ -hypernuclei in the Hyperspherical Harmonics Expansion approach

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Introduction

Hypernuclei are the subatomic systems with a remarkable strangeness degree of freedom as compared with conventional nuclei. Such exotic nuclei are formed when one or more strange exotic hyperon(s) (like Λ, Σ, Ξ etc) are injected into an atomic nucleus so as to replace one or more nucleon(s) of the nucleus and the injected particles are coupled to the nuclear core[1]. Hypernuclei provide a unique source of information on hyperon-nucleon, hyperon-hyperon interactions and the behaviour of strange particles in baryonic matter which opened an avenue for many significant phenomenon in particle physics and nuclear-astronomy[2].

Nakazawa et al., 2015[3] first reported the bound state of the p-state $\Xi^- - ^{14}\text{N}$ hypernuclei with a Ξ^- separation energy of 4.38 ± 0.25 MeV. Recently, Yoshimoto et al., 2021[4] first identified the nuclear 1s state of the hypernucleus ^{15}C having B_{Ξ^-} of IRRAWADDY (6.27 ± 0.27 MeV) and KINKA (8.0 ± 0.77 MeV) events. Some theoretical investigations were done to estimate whether the Ξ^- occupies on 1s state or 1p state in KISO event, such as relativistic mean-field (RMF) model, Skyrme-Hartree-Fock (SHF) model[5] and quark mean-field (QMF) model[6].

The main objective of the present work is to study the ground state in 1s state of ^{15}C and we will predict $^{16}_{\Xi^- \Xi^-}\text{C}$ hypernuclei. We will calculate the one- and two- Ξ^- separation energies and some relevant geometrical observables of these systems. In our two-body model, we will first solve Schrödinger equation

numerically to reproduce the single Ξ^- separation energy of ^{15}C hypernuclei by adjusting the parameter of the effective core-hyperon potential. We use the Ξ -nucleus potential together with $\Xi\Xi$ potential to predict the ground state energy of the $^{16}_{\Xi^- \Xi^-}\text{C}$ hypernuclei. For the three-body model calculation of double- Ξ system, we employ the hyperspherical harmonics expansion (HHE) method.

Theoretical formalism

In our HHE approach[7], we label the relatively heavier core nucleus ^{14}C as the particle “i” and two valence Ξ -hyperons as particles “j” and “k” form the interacting pair are:

$$\left. \begin{aligned} \vec{\xi}_i &= a_{jk}(\vec{r}_j - \vec{r}_k) \\ \vec{\chi}_i &= a_{(jk)i} \left(\vec{r}_i - \frac{m_j \vec{r}_j + m_k \vec{r}_k}{m_j + m_k} \right) \\ \vec{R} &= \sum_{i=1}^3 \frac{m_i \vec{r}_i}{M} \end{aligned} \right\} \quad (1)$$

where $\rho = \sqrt{\xi_i^2 + \chi_i^2}$; $\phi_i = \tan^{-1}(\chi_i/\xi_i)$. The motion of the three-body system in relative coordinates can be described by the Schrödinger’s equation

$$\left[-\frac{\hbar^2}{2\mu} \left\{ \frac{\partial^2}{\partial \rho^2} + \frac{5}{\rho} \frac{\partial}{\partial \rho} + \frac{\hat{K}^2(\Omega_i)}{\rho^2} \right\} \right] \psi(\rho, \Omega_i) + [V(\rho, \Omega_i) - E] \psi(\rho, \Omega_i) = 0 \quad (2)$$

where $V(\rho, \Omega_i)$ is the total interaction potential in the partition i and $\hat{K}^2(\Omega_i)$ is the square of hyper angular momentum operator satisfying the eigenvalue equation.

Choice of potential

For the core- Ξ subsystem and $\Xi\Xi$ pair, we used a two-range Gaussian Isle-type potential and Yukawa-type potential [8] of the form re-

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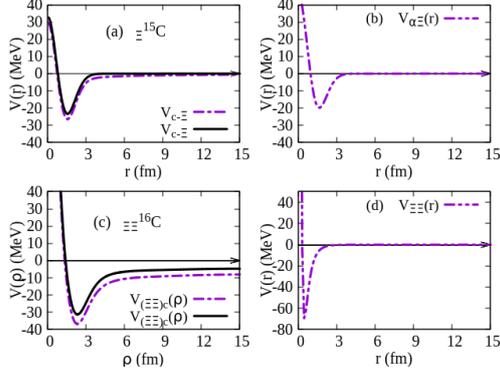


FIG. 1: Representative plot of effective core- Ξ potential components with radial distance are (a) with coulomb and without coulomb (solid line) interaction, (b) $\alpha - \Xi$ (d) $\Xi\Xi$ pair and (c) for three-body effective potential with hyper-radial distance ρ .

spectively

$$V_{c-\Xi}(r) = 450.4 \exp(-r/1.269)^2 - 404.9 \exp(-r/1.41)^2 \quad (3)$$

$$V_{\Xi\Xi}(r) = (-155.0 \exp(-1.75r) + 490.0 \exp(-5.6r))/r \quad (4)$$

A Coulomb interaction is added for the charged pair of interacting particles

$$V_{ij}^{(C)}(r_{ij}) = \begin{cases} 1.44 \left(\frac{Z_i Z_j}{2R_c} \right) \left(3 - \frac{r_{ij}^2}{R_c^2} \right) ; r_{ij} \leq R_c \\ 1.44 \left(\frac{Z_i Z_j}{r_{ij}} \right) ; r_{ij} > R_c \end{cases} \quad (5)$$

Results and Discussions

In this work, we first investigate the 1s bound state of $^{15}_{\Xi^-}\text{C}$ and predict double- Ξ hypernucleus $^{16}_{\Xi\Xi^-}\text{C}$ using HHE method which is an essentially exact method, where calculations can be carried out up to any desired precision by gradually increasing the basis.

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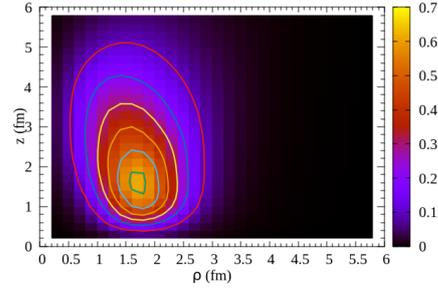


FIG. 2: Contour plot reflects probability density distribution of nucleons and Ξ^- hyperons in the ground state of $^{16}_{\Xi\Xi^-}\text{C}$.

TABLE I: (I) Parameters of the attractive component of the core- Ξ^- potential and B_{Ξ^-} separation energy of $^{15}_{\Xi^-}\text{C}$. (II) Double- Ξ^- separation energy and r.m.s raddi of $^{16}_{\Xi\Xi^-}\text{C}$. The bold face parameters obtained by switch off the coulomb interaction term. Energies and potentials are in MeV and raddi are in fm unit.

(I)	V_{att}	$V_{c-\Xi^-}$ <i>Min.</i>	B_{Ξ^-} Cal.	B_{Ξ^-} Expt.
$^{15}_{\Xi^-}\text{C}$	-417.18	-26.3556	8.0012	8.0 ± 0.77
		-23.2811	5.1671	
(II)	$B_{\Xi\Xi^-}$	R_A	$R_{c\Xi^-}$	$R_{\Xi\Xi^-}$
$^{16}_{\Xi\Xi^-}\text{C}$	-20.543	2.581	2.198	2.716
	-15.280	2.585	2.229	2.741

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